

# Lecture Nanophotonics

## Plasmonics and Metamaterials Assignment

*Class dates:* Thursday 10 April 2018  
*Due date:* Thursday 17 April 2018  
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### Note:

- We will only work with SI units.
- We use the convention in which the time-dependent part of a time-harmonic variable  $f(\omega)$  is given as  $e^{-i\omega t}$ , with  $\omega$  the angular frequency. For instance,  $E(t) = Ee^{-i\omega t}$
- The notations  $\Im(a)$  and  $\Re(a)$  are used for respectively the real and imaginary part of a complex number  $a$ .
- It may help to plot graphs on a log scale (y-axis).

## 1 Kramers-Kronig relations

The Kramers-Kronig equations constitute an important relation between the real and imaginary part of any linear response function in a physical system. They can be derived using only causality, and are therefore of very general validity.

Assume a system in which a quantity  $P$  (e.g. macroscopic polarization) is related to an external quantity  $E$  (e.g. electric field) by a linear response function  $\chi$  in the frequency domain:

$$P(\omega) = \chi(\omega) E(\omega) \quad (1)$$

Then, in the time domain:

$$P(t) = \int_{-\infty}^{-\infty} \chi(t-t') E(t') dt' \quad (2)$$

Because this is a physical system, in the time domain it must hold that:

$$\{P(t), \chi(t), E(t)\} \in \mathbb{R} \quad \text{for all } t. \quad (3)$$

Furthermore, causality implies that

$$\chi(t) = 0 \quad \text{for } t \leq 0 \quad (4)$$

because there can not be a response before there is a cause.

- a) Write  $\chi(\omega)$  in terms of  $\chi(t)$  and split in its real and imaginary part  $\chi'(\omega)$  and  $\chi''(\omega)$ .

- b) Now we will use a mathematical trick to simplify the previous result. As with any function,  $\chi(t)$  can be expressed in its even and odd parts  $\chi_e(t)$  and  $\chi_o(t)$  as

$$\chi(t) = \chi_e(t) + \chi_o(t) \quad (5)$$

Use causality to express  $\chi_e(t)$  in terms of  $\chi_o(t)$ .

*Tip: if you find this difficult, imagine  $\chi(t)$  is just a unit step function, write down  $\chi_e(t)$  and  $\chi_o(t)$  and see what the relation between the two is.*

- c) Use Eq. (5), the results from questions 1 a) and 1 b) and the convolution theorem to show that

$$\chi'(\omega) = [i\chi''(\omega)] \star \left[ \int_{-\infty}^{\infty} \text{sgn}(t)e^{i\omega t} dt \right] \quad (6)$$

where  $\star$  denotes a convolution.

- d) Show that

$$\chi'(\omega) = \int_{-\infty}^{\infty} \frac{2\chi''(\omega')}{\omega' - \omega} d\omega' \quad (7)$$

*Hint: prove the relation:*

$$\lim_{\epsilon \downarrow 0} \int_0^{\infty} e^{(i\omega - \epsilon)t} dt = \lim_{\epsilon \downarrow 0} \frac{1}{i\omega - \epsilon} e^{(i\omega - \epsilon)t} \Big|_{t=0}^{\infty} = \frac{i}{\omega} \quad (8)$$

- e) You have derived the first of the two Kramers-Kronig relations. Is it possible to have a narrow frequency window of high absorption (e.g.  $\chi''(\omega) = C_1\delta(\omega - \omega_0)$  with  $C_1 \in \mathbb{R}$ ) and no dispersion (i.e.  $\chi'(\omega)$  is constant for all  $\omega$ )?
- f) The second Kramers-Kronig relation can be derived similarly and is (you do not need to derive it):

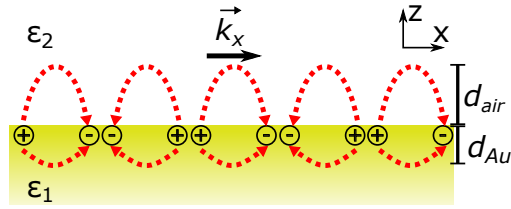
$$\chi''(\omega) = - \int_{-\infty}^{\infty} \frac{2\chi'(\omega')}{\omega' - \omega} d\omega' \quad (9)$$

Is it possible to have a window of non-zero  $\chi'(\omega)$  (e.g.  $\chi'(\omega) = C_1(\Theta(\omega - \omega_1) - \Theta(\omega - \omega_2))$  with  $C_1 \in \mathbb{R}$ ,  $\omega_1 < \omega_2$  and  $\Theta(x)$  the Heaviside step function) and no absorption?

## 2 Surface plasmon dispersion relation

The dispersion relation of a Surface Plasmon Polariton (SPP) at a metal-insulator interface is given by:

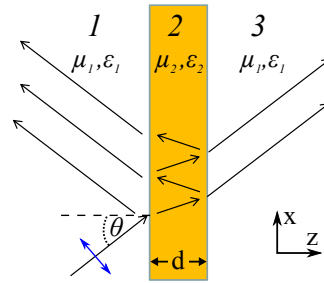
$$k_x^2 = \frac{\epsilon_1 \epsilon_2}{\epsilon_1 + \epsilon_2} \frac{\omega^2}{c^2} \quad (10)$$



- Plot the dispersion relation (real and imaginary part of  $k_x$ ) of a SPP at a gold-air interface as a function of  $\omega$  (0 to  $1.5 \omega_{p,\text{Au}}$ ). For the permittivity of gold, use a Drude model with the following parameters: plasma frequency  $\omega_{p,\text{Au}} = 2\pi \cdot 2165$  THz, damping frequency  $\gamma_{\text{Au}} = 2\pi \cdot 15.9$  THz.
- Plot the phase velocity and the group velocity of the SPP as a function of  $\omega$  and compare both to the speed of light. What happens to the velocities when  $\Re(k_x(\omega))$  reaches its maximum, and at  $\omega = \omega_{p,\text{Au}}$ ?
- The propagation length  $l$  is related to the wavevector  $k_x$  in the following way:  $l = \frac{1}{\Im(2 \cdot k_x)}$ . Plot  $l$  as a function of  $\omega$  and comment on special features and their origin.
- How deep does the electric field reach into the air medium and how deep into the gold medium? Plot  $d_{\text{air}}$  and  $d_{\text{Au}}$  as a function of  $\omega$  and explain.

### 3 Negative refraction makes a perfect lens

In this exercise, we take a look at materials with simultaneously negative permittivity  $\epsilon$  and permeability  $\mu$ . These so-called *metamaterials* (MMs) gained large attention due to their unusual electromagnetic properties, such as negative index of refraction and perfect imaging. Here, we want to consider the case of a slab of a MM (such as proposed by Veselago and Pendry) made of  $\epsilon = -1$  and  $\mu = -1$  material with a thickness  $d$  and understand what makes it a *perfect lens*. One of the most provoking claims of Veselago and Pendry is that waves experience a negative phase delay inside of such media and evanescent waves get exponentially amplified. As we will see, some of these intriguing optical properties are already evident from the transmission properties.



- Derive the formula of the total amplitude transmission  $t$  of a plane wave incident from the left side on a slab with optical properties  $\epsilon_2, \mu_2$  and thickness  $d$ , surrounded by a material with  $\epsilon_1$  and  $\mu_1$  and express it in terms of the transmission and reflection coefficients of the two interfaces of the slab.

*Hint 1: Consider all light paths contributing to the total transmission, including multiple reflected waves inside of the slab.*

*Hint 2: Use the infinite geometric series formula:*

$$\sum_{n=0}^{\infty} p^n = \frac{1}{1-p}. \quad (11)$$

- b) Consider a dielectric slab (a so-called 'etalon') with  $d = 400$  nm,  $\epsilon_2 = 10$  and  $\mu_2 = 1$  surrounded by air. Apply the derived formula to plot the intensity transmission  $|t|^2$  through this slab as a function of frequency  $\omega$  in the visible regime. Assume p-polarized light, which is incident at  $\theta = 0^\circ$  in the x-z plane ( $k_{y,1} = 0$ ). What do you see? What is the physical origin of the fluctuations?  
*Hint 1: The correct solution of a) is:*

$$t = \frac{t_{23}t_{12}e^{id \cdot k_{z,2}}}{1 - r_{21}r_{23}e^{2i \cdot d \cdot k_{z,2}}}. \quad (12)$$

*Hint 2: The Fresnel transmission and reflection coefficients for the interface between material 1 and 2 for p-polarized light are:*

$$t_{12}^p = \frac{2\epsilon_2 k_{z,1}}{\epsilon_2 k_{z,1} + \epsilon_1 k_{z,2}} \sqrt{\frac{\mu_2 \epsilon_1}{\mu_1 \epsilon_2}}, \quad r_{21}^p = \frac{\epsilon_2 k_{z,1} - \epsilon_1 k_{z,2}}{\epsilon_2 k_{z,1} + \epsilon_1 k_{z,2}}, \quad (13)$$

where  $k_{z,i} = \sqrt{k_i^2 - k_{\parallel}^2}$  and  $k_i = \frac{\omega}{c} \sqrt{\epsilon_i \mu_i}$  with material indices  $i = 1, 2$ .

- c) Use the same equation as in b) to plot  $|t|^2$  and  $\arg(t)$  as a function of the normalized parallel wavevector in the interval  $\{0 \leq k_{\parallel}/k_1 \leq 2\}$  at  $\omega = 2\pi \cdot 375$  THz. Assume that all materials are air ( $\epsilon_1 = \epsilon_2 = 1$  and  $\mu_1 = \mu_2 = 1$ ) and  $d = 50$  nm. What do you see? What does  $k_{\parallel} > 1$  mean for the incidence angle and  $k_z$ ? What is the physical significance of  $\arg(t)$ ?
- d) Describe qualitatively how the amplitude of propagating ( $\omega^2/c^2 > k_{\parallel}^2$ ) and evanescent ( $\omega^2/c^2 < k_{\parallel}^2$ ) waves changes as a function of distance if the slab is made of air (*hint: use equation 12*). And what happens in the case of negative index MM ( $\epsilon_2 = -1$ ,  $\mu_2 = -1$ )? What are the directions of  $k_{z,2}$ , Poynting vector, phase advancement and energy flux in the negative index MM?
- e) Make the same plots as in c) at  $\omega = 2\pi \cdot 375$  THz for a slab of MM with  $\epsilon_2 = -1$  and  $\mu_2 = -1$  and  $d = 50$  nm. Pay attention to the orientation of  $k_{z,2}$ . What is the surprising part about these plots as compared to c)?